

The Influence of Peptizers on the Efficiency of a Plodder in SOV/72-59-7-3/29
Manufacturing Faience Materials and Kaolin-bentonite-suspensions

plodder the current consumption being measured by means of a self-recording wattmeter of the system D-333. The test results are represented in figure 1 till 5 and subsequently explained. The influence of the peptization on the processing conditions is evident both from the former investigations of the authors of this paper and F. D. Ovcharenko (Footnote 4) and from the investigations of G. V. Kukolev and Ya. M. Syrkin (Footnote 5). The maximum output of the plodder is attained by constant number of revolutions of the worm shaft at the optimum suspension viscosity. The test results agree with the former studies of the authors of this paper and G. V. Kukolev and A. N. Korol' (Footnotes 6 and 7). Conclusions. Bentonite additions to kaolin suspensions and faience materials increase the output of plodders considerably and lower the current consumption. An even greater effect is caused by the addition of some peptizers. There are 5 figures, 2 tables, and 11 references, 9 of which are Soviet

Card 2/2

PITAK, N. V.

Cand Tech Sci - (diss) "Effect of bentonite on the properties of clayey suspensions, faience gross, mass and finished articles." Khar'kov, 1961. 20 pp; (Ministry of Railways USSR, Khar'kov Inst of Railroad Transport Engineers imeni S. M. Kirov; 1-20 copies; free; list of author's works at text's end (10 entries); (KL, 6-61 sup, 223)

KUKOLOV, G.V.; STRELETS, V.M.; PITAK, N.V.; AMERIKOVA, T.A.

Sectional nozzles for the continuous pouring boiling steel. Ogneporov
25 no.8:352-356 '60. (MIRA 13:9)

1. Ukrainskiy nauchno-issledovatel'skiy institut ogneporov.
(Steel--Metallurgy)

KUKOLEV, G.V., [Kukoliev, H.V.]; PITAK, N.V. [Pytak, M.V.]

Change in the properties of bentonite suspensions due to the effect of peptizers. Dop.AN URSSR no.6:772-775 '61.

(MIPA 14:6)

1. Khar'kovskiy politekhnicheskii institut. Predstavleno akademikom AN USSR F. P. Budnikovym.

(Bentonite)

(Colloids)

15(2)

V. 72-11-11

AUTHORS: Kukolev, I. V., Professor, Institute of Technical Sciences,
Pitak, B. V.

TITLE: Utilization of Bentonite in the Production of Faience Products
(Ispol'zovaniye bentonita dlya proizvedstva fayansovykh izdeliy)

PERIODICAL: Steklo i Keramika, 1964, No. 1, pp. 25-26 (USSR)

ABSTRACT: The addition of bentonite increases the plasticity of the pastes and increases the solubility and water stability of the semifinished products in the air-dry state, whereby a transition to a single burning of the products is rendered possible. Moreover, bentonites as fillers can reduce the burning temperature of ceramic products, as may be seen from papers published by G. P. Filinsev, M. A. Bezbardov, B. G. Puzanov, and V. P. Shvayko (Ref. 1). In the field of the use of bentonite in faience pastes there exist only a few papers by B. A. Bezbardov, Ye. F. Golubtsov, T. A. Nisova, and A. Shadrin, and I. Ya. Iiven' (Ref. 2). It may be seen therefrom that bentonites can be used for the production of faience tiles and slabs. The authors of this article made experiments in this field with the chemical composition of the raw materials used hereinafter shown in table 1. The viscosity of clay was determined by means of the viscosimeter V-1, figures 1, 2.

Card 1 2

Utilization of Bentonite in the Production of Paper
Products

and 3 show the effect of the type and the quantity of additions on viscosity. Table 2 shows the compositions and essential properties of the clays and products. In connection with the performance of these experiments the authors of this article refer to papers published by G. V. Mukolev, Ya. M. Syrkin (Ref 2), I. A. Kryukov, and N. A. Komarov (Ref 4). Figure 4 shows the dependence of the ability to retain water on pressure and figure 5 the volumetric changes of the samples in drying at 110°. The filterability of the clays was determined by means of a vacuum filter which is described in papers published by Ya. M. Syrkin, L. N. Bernshteyn, and K. L. Kiseleva (Ref 5). The best filterability of clay occurs with an increased kaolin content and a bentonite addition as may be seen from table 3. Conclusions: Products of this clay show less shrinkage after drying and burning than products of lime with a content of chasov-yarskiy clay and exhibit a considerable water retention. The water absorption of the products can be reduced to 10% by the addition of 1% filler to the clay. (Refer to figures, tables and 5 Soviet references.)

Card 2, 2

STRELETS, V. M., PITAK, N. V.

Increasing the strength of stoppers of 140-ton steel-pouring
ladles. Ogneupory 25 no.4:171-175 '60. (MIRA 13:8)

1. Ukrainskiy nauchno-issledovatel'skiy institut ogneporov.
(Steelworks--Equipment and supplies)
(Refractory materials)

KUKOLEV, G.V.; PITAK, H.V.

Effect of the type and amount of peptizators on the modification
of structural and mechanical properties of a bentonite suspension.
Izv. vys. ucheb. zav.; khim. i khim. tekhn. 2 no.2:244-246 '59.
(MIRA 12:9)

1. Khar'kovskiy politekhnicheskii institut imeni V.I. Lenina.
Kafedra keramiki, stekla, ogneuporov i emalirovaniya.
(Bentonite)

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S/131/60/000/04/05/015
B015/B008

18 4000 15.2200

AUTHORS: Strelets, V M., Pitak, N V.

TITLE: Increasing the Stability of Stoppers of 140 t Steel-casting Ladles

PERIODICAL: Ogneuporny. 1960, No. 4, pp. 171-172

TEXT: In the paper under review the authors describe the function of the chamotte pipes SP-8-2, SP-8-4 and the chamotte stoppers SP-13-1 of the Zaporozhskiy ogneupornyy zavod (Zaporozh'ye Works for Refractories), the quartz-kaolin pipes SP-8 of the Prosyanyovskiy kaolinovy kombinat (Prosyanyaya Kaolin Kombinat), magnesite sleeve bricks of the Chasov-Yarskiy kombinat ogneupornykh izdeliy (Chasov-Yar Kombinat for Refractories) and sleeve bricks of the Konstantinovskiy ogneupornyy zavod "Krasnyy Oktyabr'" (Konstantinovka Works for Refractories "Krasnyy Oktyabr'") I I. Druzhinin, Yu Z Babaskin, and A N Slin'ko participated in the experiments. The physicochemical properties of the materials used are mentioned in table 1. The pipes are corroded most by slag (Fig. 1). Examples of the wear of the pipe seams and the sleeve bricks are shown in Figs. 2 and 3 and the varied insulation of the stopper rods in Fig. 4. Mortar of varied composition was tested in the experiments (Table 2) in order to

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Increasing the Stability of Stoppers of
140 t Steel-casting Ladles

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eliminate the corrosion of the pipe seams. The authors in conclusion underline that the amount of slag in the ladle constitutes one of the main factors for the corrosion of the stopper pipes. The tearing-off of the spherical part of the stopper, caused by the formation of a crust between sleeve brick and stopper, can be eliminated by a graphite covering. The corrosion of the pipe seams may be reduced by using quality mortar for the insulation of the stoppers. A highly aluminous coating of the stopper pipes eliminates their wear. There are 4 figures, 2 tables, and 9 references, 8 of which are Soviet.

ASSOCIATION Ukrainakly nauchno issledovatel'skiy institut ogneporov (Ukrainian Scientific Research Institute of Refractories)

PITAK, N.Y.

Nozzles for the semicontinuous casting of stainless steel
under synthetic slag. Ogroupory 1964. 32. 1-17. 1964.
(MI 1964)

1. Ukrainskiy nauchno-issledovatel'skiy institut obrabotki met.

5(1,2)

AUTHORS:

Kukolev, G. V., Pitak, N. V.

SOV/153-2-2-19/31

TITLE:

Change of the Structural and Mechanical Properties of the Bentonite Suspension in Its Dependence upon Kind and Quantity of Peptizers (Izmeneniye strukturno-mekhanicheskikh svoystv bentonitovoy suspensii v zavisimosti od roda i kolichestv peptizatelev)

PERIODICAL:

Izvestiya vyzhikh uchebnykh zavedeniy. Khimiya i khimicheskaya tekhnologiya, 1959, Vol 2, Nr 2, pp 244 - 246 (USSR)

ABSTRACT:

The questions of the influence of various electrolytes and surface-active substances on the suspension mentioned in the title (Refs 1-3) have not been investigated sufficiently. Some electrolytes cause the development of coagulation structures (Ref 4). Alkaline electrolytes accelerate and harden the thixotropic structures of the bentonite suspension (Ref 5). NaOH liquefies such a suspension, but only in the absence of CO₂ of the air.

Pyshevskiy bentonite was examined (Khmel'nik area UkrSSR). According to Ovcharenko (Ref 7), the compound water here amounts to 20% and the specific surface to 736 m²/g. A rotary viscosimeter of the type HV - 8 (Ref 8) served for its determination at room

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Change of the Structural and Mechanical Properties of the SOV/153-2-2-19/31
Bentonite Suspension in Its Dependence Upon Kind and Quantity of Peptizers

temperature. Peptizers were: soda, caustic soda, soda extracts of peat and straw, water glass, peat extract with water glass, vinasse of sulfite spirit and tannin. The concentration of solid substances in the system amounted to 33%. The suspension was left standing for 8-10 hours. The structure developed in that time (Ref 9), was carefully destroyed by shaking. Figure 1 and 2 show the results of the definition. From them one can see that the peptizers change the viscosity and the threshold shear stress of the bentonite suspension in a much larger field of sizes, with the same moisture content. According to their effect, peptizers may be divided into 3 groups:

1. those which reduce the solidity of the structure - liquefiers;
2. those which increase this solidity;
3. surface-active admixtures which take an intermediate position between 1. and 2. and which have little influence on the indices of the structural mechanical properties. 1. includes: water glass and water glass extract of peat. The bentonite in suspension, together with these admixtures, is peptized under the influence of alkali developed due to the hydrolysis of the alkaline electrolyte. Thus the penetration of water into the interior of the structure is favored.

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Change of the Structural and Mechanical Properties of the SOV/153-2-2-19/31
Bentonite Suspension in Its Dependence Upon Kind and Quantity of Peptizers

Colloidal silicic acid however envelops the bentonite particles or their aggregates with a protective film of sufficient strength. In this way the penetration of water mentioned before is prevented. Thus conditions are created under which the liquefying influence of the alkaline cations and OH ions on the bentonite suspension is made effective. That is how a considerable reduction of viscosity and of the threshold shear stress develops. 2. Peptizers of this group: soda, caustic soda, soda extracts of peat and straw increase the dispersion of the bentonite pools most. Thus the mechanical properties of the suspension are increased (curves 1-4, figures 1 and 2). Na^+ and OH^- ions liquify the suspension only in the absence of CO_2 (Ref 6). 3. Among them there are the surface-active organic substances: sulfite spirit, vinasse and tannin. The curves 5 and 6 show a weak effect on the reduction of the structural and mechanical properties. There are 2 figures and 10 references, 9 of which are Soviet.

ASSOCIATION:

Card 3/4

Khar'kovskiy politekhnicheskiy institut imeni V. I. Lenina; Kafedra keramiki, stekla, ogneporov i enalirovaniya (Khar'kov Polytechnic Institute imeni V. I. Lenin, Chair of Ceramics, Glass, Refractories,

Change of the Structural and Mechanical Properties of the **SOV/153-2-2-19/31**
Bentonite Suspension in Its Dependence Upon Kind and Quantity of Peptizers

and Enamelling)

SUBMITTED: March 5, 1958

Card 4/4

AUTHORS: Kukolev G V and Pitak, N V 21-58-5-21/28

TITLE: Water Retention Properties of Kaolin, Bentonite and Faience Masses (Vodouderzhivayushchaya sposobnost' kaolina bentonitov i fayansovykh mass)

PERIODICAL: Dopovidi Akademii nauk Ukrain's'koi RSR 1958, Nr 5 pp 549-553 (USSR)

ABSTRACT: The authors describe their experiments in determining the water retention properties of various materials such as kaolin, bentonite obtained from different regions, clay and faience masses. The experiments have shown that bentonite possess the greatest retainability. These experiments were carried out under various pressures and their results are represented by graphs showing the amount of water (V) plotted versus the pressure applied (P). The P-V curves of kaolin suspensions and faience masses show horizontal or sloped sections in certain ranges of pressure. Peptizers increase the water retention properties of pure kaolin and bentonite suspension at all pressures, as well as their mixtures, and of faience masses at low pressure (up to 50 kg/cm²). It is shown that the P-V curves are described by empirical formulas. Equations for the calculation of the quantity Δ (being a derivative

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21-58-5-21/28

Water Retention Properties of Kaolin Bentonite and Faience Masses

of pressure with respect of water amount (V) are derived from these formulas. Given is the dependence of Δl on the values of P and the kind of the peptizers used. There are 4 graphs, 1 table and 9 Soviet references.

ASSOCIATION: Khar'kovskiy politekhnicheskiy institut (Khar'kov Polytechnic Institute)

PRESENTED: By Member of the AS UkrSSR P. P. Budnikov

SUBMITTED: October 24, 1957

NOTE: Russian title and Russian names of individuals and institutions appearing in this article have been used in the transliteration.

Card 2/2

8/13/82 000/000, 000, 000
2117/2101

Authors: Strelits, V. M., Pitak, N. V., Kulik, N. I., Bogachov, N. S.

Title: Laboratory investigations of the technology of zircon products

Journal: Ogneupory, no. 6, 1968, 183-188

Summary: The influence of the following factors on the physico-chemical properties of zircon products was studied: grain composition, molding pressure, burning temperature, admixtures of clay, raw zircon concentrate (ZrO₂ 210-47, ISMZU 200-47), and raw non-ferrous zircon (ZrO₂ 400-94, ISMZU 400-94), the object being to establish optimum masses and working standards for the production of proportioning ladles for use in continuous steel-casting foundries. The lowest apparent porosity and the highest weight by volume were determined after drying (at 120°C) of samples made up of 1.5-0.5 mm grains (50%) and of < 0.085 mm grains (50%), and after burning (at 1550°C for 2 hrs) of samples made up of 1.5-0.5 mm grains (50%) and of < 0.085 mm grains (70%). A pressure of 500 kg/cm² was found sufficient for the production of dosing cups, as an increase in

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Laboratory investigations of ...

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1117/5101

pressure from 500 to 1850 kg/cm² reduces the apparent porosity of
 15-100 only. The fine grain size (< 0.008 mm) of burned zircon
 be replaced by the same grade of raw zircon. An increase of the burning
 temperature from 1550 to 1600°C raises the linear shrinkage of
 100% and the compressive strength from 40-50 to 60-100 kg/cm².
 100% clay improved the plasticity and made molding easier. High-
 strength products (1000 kg/cm²) were obtained at lower temperature
 1500-1550°C. Raw zircon and zircon concentrate may be used for smaller
 100% products, which must be burned at < 1550°C to avoid swelling.
 addition of clay reduces the temperature of sample destruction from
 1000 kg/cm² by 150-200°C. This temperature reduction is smaller
 100% plus of burned zircon. There are 4 figures and 2 tables.

ASSOCIATION: Укрainskij nauchno-issledovatel'skij institut keramiki
 Ukrainian Scientific Research Institute of Refractory
 Materials (Strelots, V. M., Patak, M. V., Chukovskiy,
 Komsinat ognepornyykh izdeliy (Chusov Yar Combine of
 Refractory Products) (Kulik, A. I., Logachev, M. S.)

Car. 2/2

S/131/60/000/008/001/003
B021/B058

AUTHORS: Kukolev, G. V., Strelets, V. M., Pitak, N. V.,
Amerikova, T. A.

TITLE: Compound Pouring Ladle Nozzle Lining for the Casting of
Rimmed Steel in Installations for Continuous Steel Casting

PERIODICAL: Ogneupory, 1960, No. 8, pp. 352-356

TEXT: It was the authors' task to elaborate a ladle nozzle lining, which undergoes only slight wash-out, is not clogged by metal, and warrants a satisfactory jet without spattering or eddies. Highly aluminous zirconium- and magnesite inserts for the compound pouring ladle nozzle lining were produced at the Opytnyy zavod (Experimental Plant) of the UNIIO (Ukrainskiy nauchno-issledovatel'skiy institut ogneuporov - Ukrainian Scientific Research Institute of Refractories). The pouring ladle nozzle linings were produced at the Chasov Yarskiy kombinat ogneupornykh izdeliy (Chasov Yar Kombinat of Refractories), the working processes having been previously elaborated at the Experimental Plant of the Ukrainian Scientific Research Institute of Refractories. Technical alumina of type П 1 (G1) and Chasov

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Compound Pouring Ladle Nozzle Lining for the
Casting of Rimmed Steel in Installations for
Continuous Steel Casting

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B021/B058

Yar clay 41 (Ch1) were used for the production of highly aluminous¹⁵ inserts. Zirconium inserts were produced from finely ground zirconium with a ZrO_2 content of 69%. Chamotte pouring ladle nozzle linings were produced at the Experimental Plant of the Ukrainian Scientific Research Institute of Refractories from a mass containing 40% chamotte from Chasov Yar clay 41 (Ch1), 40% Chasov Yar clay 41 (Ch1) and 20% foundry coke. The highly aluminous and magnesite inserts, as well as chamotte pouring ladle nozzle linings were pressed in the "Tagilets" friction press. A press mold (Fig. 1) was used at the Chasov Yar Kombinat. A total view of the two parts of the compound pouring ladle nozzle lining is shown in Fig. 2. The inserts and linings were fired in periodic furnaces. The firing curves are shown in Fig. 3 and the properties of the fired products are tabulated. The compound linings were tested at the Stalinskiy metallurgicheskiy zavod (Stalino Metallurgical Plant) and the zavod "Krasnoye Sormovo" ("Krasnoye Sormovo" Plant) during the casting of rimmed steel. The experiments were conducted by collaborators of the Ukrainian Scientific Research Institute of Refractories, the Ukrniimetallov (Ukrainskiy nauchno-issledovatel'skiy

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Compound Pouring Ladle Nozzle Lining for the
Casting of Rimmed Steel in Installations for
Continuous Steel Casting

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institut metallov - Ukrainian Scientific Research Institute of Metals),
the TsNIICHM (Tsentral'nyy nauchno-issledovatel'skiy institut chernoy
metallurgii - Central Scientific Research Institute of Ferrous Metallurgy).
the Stalino Metallurgical Plant and the "Krasnoye Sormovo" Plant. Fig. 4
shows highly aluminous inserts after their use in 50 t pouring ladles.
They were tested at the "Krasnoye Sormovo" plant with apertures of 30 mm
diameter. The aperture of the insert was washed out by 1-2 mm in diameter
when casting rimmed steel of type 3кп (3kp). The wear amounts to 4-6 mm
when casting weld steel of type CB 08A (Sv08A), which is explained by its
higher content of iron oxides. The authors state in conclusion that the
production technology of compound nozzle linings was elaborated for con-
tinuous rimmed-steel casting. The compound lining consists of a porous
chamotte pouring ladle nozzle as a carrying part, and a highly aluminous
magnesite- or zirconium insert as working part. The highly aluminous
inserts showed the best wear resistance during tests. There are 4 figures,
1 table, and 5 references: 1 Soviet, 2 British, and 2 US.

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Compound Pouring Ladle Nozzle Lining for the
Casting of Rimmed Steel in Installations for
Continuous Steel Casting

S/131/60/000/008/001/003
B021/B058

ASSOCIATION: Ukrainskiy nauchno-issledovatel'skiy institut ogneporov
(Ukrainian Scientific Research Institute of Refractories) ✓

Card 4/4

15(2)

AUTHORS: Kuz'min, I. I., Strelets, V. I.

TITLE: Application of the Method of Powder Compaction for the
Basing of the

PERIODICAL: Chemistry, 1964, 11, 10-11

ABSTRACT: In the "Kuz'min" method of working together
considered of the Ironyanyaya
of fine Refractor, of the
clay these
aluminum Refractor. The
per were produced
of the semi-dry
method. The physical and chemical properties of stopper
are listed in table 1, their wear may be seen from table 2.
In figure 1 and the
to their use according to their use. The chemical
is indicated in table 3. The microstructure of
buckling is given in
figures 2 and 3. It is concluded
that the is brought

Card 1/2

Application and Variation of Phase Composition of the Stopper Bushing of Casting Ladles in Continuous Steel Casting

SOV/131-59-12-6/15

about mainly by the action of the slag and of the molten metal. The greatest stability is found with bushings of high alumina content. It is considered interesting to investigate the possibility of prolonging life of fire clay lining of the casting ladle and stoppers by the addition of grog. The possibility of using covers for casting ladles should be investigated in order to be able to cast with a minimum slag cover. There are 4 figures, 3 tables, and 9 references, 8 of which are Soviet.

ASSOCIATION:

Ukrainskiy nauchno-issledovatel'skiy institut ogneniy rovy
(Ukrainian Scientific Research Institute of Refractories)

15(2)

AUTHORS:

Strelets, V. M., Pitak, N. V.

S/131/60/000/01/009/017
B015/B001

TITLE:

Experiments on the Use of Sleeve Bricks for Continuous Steel Casting

PERIODICAL:

Ogneupory, 1960, Nr 1, pp 30 - 32 (USSR)

ABSTRACT:

In this paper, the authors describe experiments with sleeve bricks with different sleeves (Fig 1). N. P. Mayorov, N. S. Agazor'yants, A. V. Khrbkov, A. M. Makushin, L.B. Shenderov, V. G. Barsukov, and Z. D. Abuladze participated in the experiments. Table 1 shows the chemical composition of the sleeve bricks and the sleeves. The casting conditions of steel and the wear of the sleeve bricks in a plant for continuous steel casting are given in table 2. Figure 2 shows a biceramic sleeve brick with a layer of high alumina content after use. In conclusion, the authors mention that unburnt sleeve bricks with a magnesite layer show a higher wear resistance than those with a clay-graphite layer. Sleeves of highly refractory materials showed the highest durability. There are 2 figures and 2 tables.

Card 1/2

STRELETS, V.M.; PITAK, N.V.

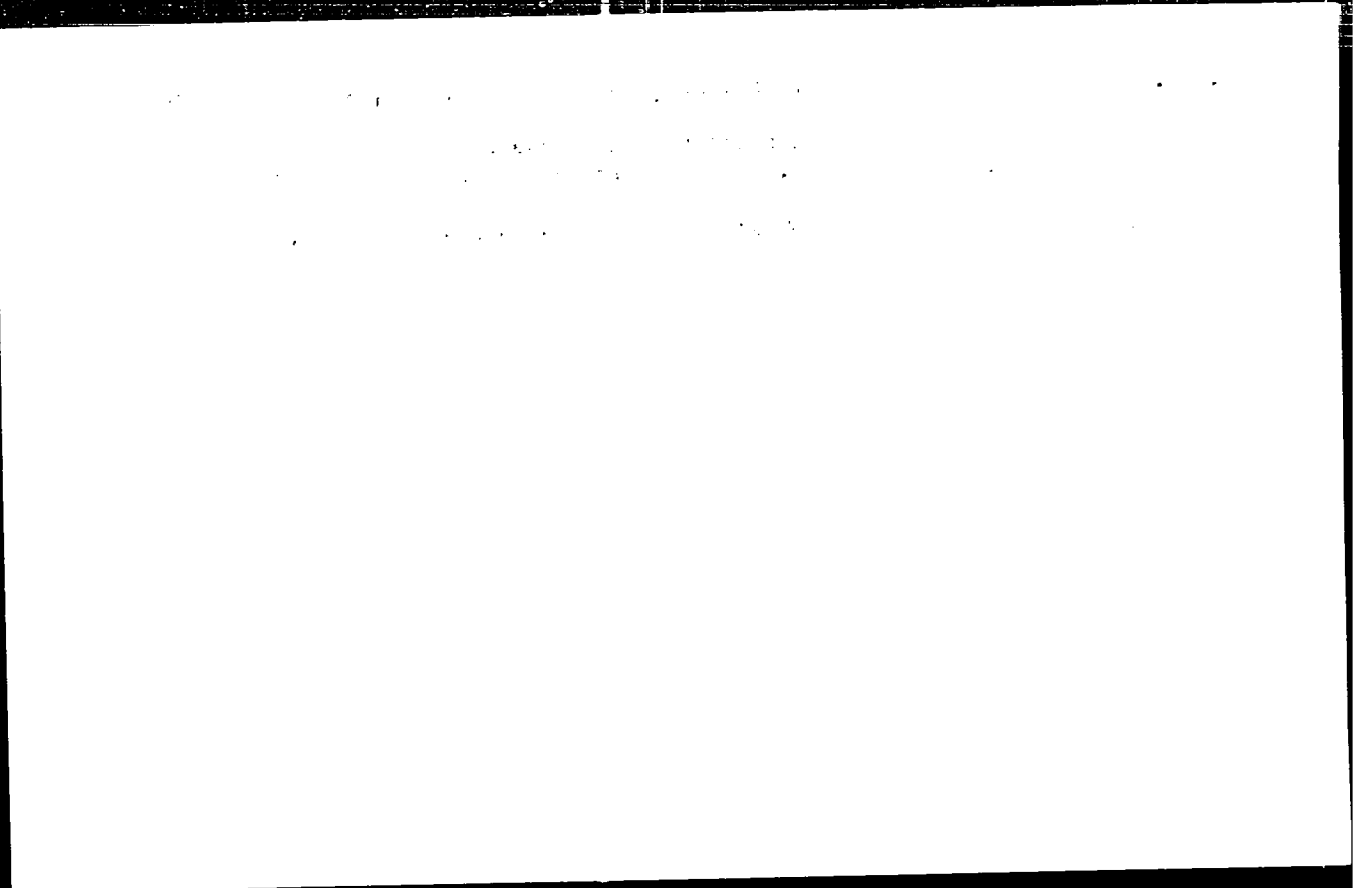
Service characteristics of stoppers during the continuous pouring of
steel. Ogneupory 25 no.2:64-69 '60. (MIRA 13:10)

1. Ukrainskiy nauchno-issledovatel'skiy institut ogneuporov.
(Refractory materials) (Steel--Metallurgy)

STRELETZ, V. A. ; PIVAR, N. V. ; KULIK, A. I. ; LOGACHEV, M. S.

Laboratory investigation on the technology of zirconium
Grouping: 15.6.283.288. 62.

Исследования научно-исследовательский институт
1 - группа атомной энергии. Часов-варский комбинат
Искандер (for Kulik Logachev).
(Zirconium Testing)



PITAMIC, Tomo, doc. d.

Treatment of rheumatoid arthritis in childhood. Reumatizam 12
no.3:83-88 '65

1. Klinika za dječje bolesti Salata medicinskog fakulteta u
Zagrebu.

PITAMIC, T., doc. dr.; GRGIC, M., dr.; SPIDLA, M., dr.

Sedimentation of erythrocytes in the treatment of rheumatic fever.
Reumatizam 12 no.1:7-11 '65

1. Klinika za dječje bolesti Salata Medicinskog fakulteta u
Zagrebu.

opisni prikaz, Lea, prof. dr.; LITAMIC, Tomo, doc. dr.; GRGIC, Miljenko, dr.;
MARTIĆ, Marijana, dr.

Changes in cardiac findings in the course of recurrent rheumat.
fever. Reumatizam 12 no.6:211-216 1965.

. KAKA ZA Njele no esti Saraja Medicinskog fakulteta .
AR 0001

GBERHOFFER-SIK, Tea, prof. dr. PITAMIC, Tomo, doc. dr.; GRGIĆ, Miljanka, dr.;
MARK, Bruno, doc. dr.; BUNETA, Dragutin, dr.

Changes in the development of carpal bones in congenital heart
defects. Reumatizam 12 no.2:41-47 '65

I. Klinika za dječje bolesti Salata i Zavod za radiologiju Medi-
cinskog fakulteta u Zagrebu.



MAMLIN, O.A., insh.; PITANOVA, M.S., insh.

Span for a new bridge in Lon' grad. Transp.stroi. 14
no.12:29-32 D '64 (MIRA 19:1)

BOROVIKOV, V.N., inzh.; MAMLIN, G.A., inzh.; PITANOVA, N.S., inzh., REUT, Z.V.,
inzh.

Preparing welded, box elements for span structures. Transp. street 14
no. 7. 23-26 JI '64. (MIRA 18.1)

1. 12953-63
ASD/ESD-3

EMD(k)/EWP(q)/EWT(m)/EWT(l)/BDS/T-2/EEC(b)-2/ES(t)-2 AFPTC/
Pa-4/Pa-4 IJP(C)/JD 8/109/63/008/004/029/030 75

AUTHORS: Logunov, L. A., and Pitanov, V. S.

TITLE: Volt-ampere characteristics of tunnel diodes made of gallium arsenide

PUBLICATION: Radiotekhnika i elektronika, v. 8, no. 4, 1963, 723-725

TEXT: The authors report on their research into the inverted and direct branches of the volt-ampere characteristic of pn-junctions obtained with gallium arsenide of the p-type, alloyed with zinc [in a concentration of $(4-10) \cdot 10^{19} \text{cm}^{-3}$]. The drop in voltage r_s in the series resistance of the diode was taken into account in plotting the characteristics of the pn-junction. The measurements were taken while passing, through the diode, currents of from ~ 100 to 300 ma, for a duration of $\sim 1 \mu$ -sec. The voltage drop was measured in the diode, while the impulse current was determined on the basis of the amount of voltage drop when using a resistance of 10 ohms connected in series with the diode. On the basis of two measurements of the voltage at various high current values, the value of r_s could thus be determined. Two graphs accompany the article. One shows the relationship between the logarithm of conductivity of the pn-junction and the respective pn-junction voltage. This relationship is very close to rectilinear. The other graph shows the relationship between the logarithm of the direct current and the

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Volt-ampere characteristics

shift in the pn-junction. In the negative sector of the characteristic, this relationship is closely approximated by an exponential curve.

SUBMITTED: December 18, 1962

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LOGUNOV, L.A.; PITANOV, V.S.

Voltampere characteristics of gallium arsenide. Radiotekh. i
elektron. 8 no.4:723-725 Ap '63. (MIRA 16:4)
(Tunnel diodes)

Pituleva R. A.

*Influence of Camphor on the Electrode Potential of the Dropping Amalgam Electrode, A. G. Stromberg and R. A. Pituleva (*Doklady Akad. Nauk S.S.S.R.*, 1953, 89, (8), 1071-1074; errata, 1954, 97, (5), 866).—[In Russian]. The potentials of dropping electrodes of Pb, Sn, Cd, and Zn amalgams in 0.005M soln. of ions of the metals in various indifferent electrolytes (0.1N-KNO₃ or 0.1N-HNO₃; 0.2N-HCl or 0.1N-KOH; 0.1N-KCl; 0.1N-KCl or 0.5N-KOH, resp.) shifted in the positive direction on saturating the soln. with camphor: there was no change with Bi amalgam and 0.5N-HCl, whilst with Cu amalgam and 0.1N-H₂SO₄ there was a slight shift in the negative direction. Reducing the drop period from 2 to 1 sec. (thus increasing the rate of formation of a new electrode surface from 0.013 to 0.026 cm.²/sec.) in experiments with Cu, Bi, and Pb amalgams led to shifts of -0.003, 0, and +0.010 V., resp. In the case of Cd amalgam, shifts of +0.120, +0.041, and +0.003 V. were produced when the soln. was saturated, $\frac{1}{2}$ -saturated, and $\frac{1}{4}$ -saturated, resp., with camphor. Increasing the Cd content in the amalgam (0.0013M) by 10x and 100x reduced the shift from 0.112 V. to 0.043 and 0.003 V., resp. Camphor did not affect the equilibrium potentials of stationary amalgam electrodes. If camphor is absent, the potentials of dropping and stationary electrodes are identical. All these observations can be explained in terms of the theory that the electrode processes are governed by exchange currents.—G. V. E. T.

MG
3

10

PITAS V.

Strojrenstvi (Machinery)
Vol. 8, Nr. 1, 21. 1958

Pitas V.: Trichloroethylene¹ filtering installa-
tions.

Strojrenstvi, Vol. 8, No. 1, 1958, p. 53-54

The article deals with a novel type of filter to be used in trichloroethylene installations. The filter has many advantages over conventional filters, centrifugal cleaning and evaporating devices. The filter is very economical and cheap in operation requiring only 4% soap solution.

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PITAELEVA, R.A.

USSR.

Influence of campher on electrode potential of dropping amalgam electrode. A. G. Strumborg and R. A. Pitaeleva (A. M. Gor'ki Ural State Univ., Sverdlovsk). *Doklady Akad. Nauk S.S.S.R.* 99, 1071-4(1968).—Electrode potentials (E) of dropping or stationary electrodes of Cu, Bi, Pb, Sn, Cd, and Zn amalgams were measured in 0.005M solns. of salts of corresponding metals in aq. HCl, H₂SO₄, HNO₃, KOH, KCl, and KNO₃ solns., both in the presence or absence of campher in the electrolyte. Campher caused shifts in E : in the case of Cu, the shift was in a neg. direction; Bi had no influence; and with all other metals the shift was in a pos. direction. Increase in the dropping rate in campherated solns. shifted E in the case of Cu amalgam in a neg. direction and in the case of Pb amalgam in a pos. direction, and had no influence in the case of Bi amalgam. The magnitude of the shifts increased with increase in campher concn. With increase of CO concn. in the amalgam, the deviations caused by campher decreased. With stationary amalgam electrodes, E in campherated solns. were not different from E without campher. E of dropping amalgam electrodes were equal to that of the stationary electrodes if campher was not present. These observations agree with predictions which follow from the theory that the E of the amalgam electrodes are governed by exchange currents. Campher retards the exchange currents and thus shifts E in the direction dependent on the relative position of E of the metal with respect to the zero point of the dropping electrode. With faster dropping, less time is available for a full buildup of the exchange currents. If enough time is allowed, as with the stationary electrode, campher has no influence on E since it forms no complexes. Without campher, the buildup of the exchange currents is fast and E of the dropping amalgam electrode become equal to that of the stationary one. Increase in concn. of the metal in Hg intensifies the exchange currents so that campher has less influence.

Amirav Ibravich

62

PITANIYE I.

42648. Obmen Veshchesty Pri Fosfornom Otravlenii. [~]obshch.11. G.P. Yeregin.
Obmen Kal'ts'ya, Fosfora I Obshchego Kolichestva Osnovaniy U Otravlennykh Fosforom
Zhivotnykh Pri Razlichnom Pitanii. Gigiena I Sanitariya, 1948, No. 12, S. 25-32.

PITAYEV, L. P.

SUBJECT: ...
 AUTHOR: ...
 TITLE: ...
 REVISIONS: ...

It is known that ... the spectrum of ... wave function ... Here n denotes the ... the hamiltonian ... by a function ... order of ϵ^n . The term ... $\int \dot{v}^2 dt + \dots$ with respect to n and consider the ... the fact that ... r_{α}^{β} ... oscillators with the ... of ϵ it is assumed that ... in the basic state ...

ARKHIPOV, R.G. [translator]; GOR'KOV, L.P. [translator]; DZYALOSHINSKIY,
I.Ye. [translator]; PITAYEVSKIY, I.P. [translator]; KHALATNIKOV,
I.M., red.; ~~BEKBERDIN~~, I.M., red.; KHOMYAKOV, A.D., tekhn.red.

[New properties of the symmetry of elementary particles.
Translated from the English] Novye svoystva simmetrii elemen-
tarnykh chastits; sbornik statei. Perevod s angliiskogo
R.G.Arkhinova i dr. Moskva, Izd-vo inostr.lit-ry, 1957. 97 p.
(MIRA 11:1)

(Particles. Elementary)

PITAYEVSKIY, L. P.

Derivation of the formula for the energy spectrum of
liquid helium - L. P. Pitayevskiy

1. RMB
2. 1/2

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ПИТАЙЕВСКИЙ

AUTHORS

Lifshits, Ye.M., Pitayevskiy, L.P.

56-2-35/47

TITLE

On the Absorption of the Second Sound in Rotating Helium II.

(O pogloshchenii vtorogo zvuka vo vrashchayushchetsya gellii II)

PERIODICAL

Zhurnal Eksperim. i Teoret. Fiziki, 1957, Vol. 33, Nr 2(8), pp 535-537 (USSR)

ABSTRACT

The authors are inclined to look upon the good qualitative agreement of theory with the experiment, which was obtained by H.E. Hall and W.F. Vinen (Proc. Roy. Soc., Vol. A, 238, pp. 204, 215) as largely due to chance. Besides, an indefinite empirical constant occurs in this calculation.

The authors computed the scattering of rotons by the vortex threads in quasiclassical approximation. For the force acting upon the unit of length of the thread the following result is obtained (the method of denotation being the same as with H.E. Hall and W.F. Vinen):

$$\vec{F} = D (\vec{v}_R - \vec{v}_L) + D' [\omega, \vec{v}_R - \vec{v}_L] / \omega, D \sim 1, 2$$

$$\sqrt{\mu} k T / p_0, D' = \chi q_R$$

CARD 1/3

56-2-35/47

On the Absorption of the Second Sound in Rotating Helium II.

Circulation round the vortex is assumed to be equal to $\kappa = 2\pi \hbar/m$. Next, two further coefficients are computed. Comparison with experimental data is, of course, possible only in that domain in which it is possible to speak of a "roton gas" and in which the range of the rotons is small in comparison with the spacings between the vortex threads. The computed values of coefficients are shown in form of a diagram. Though they agree well with the measurements carried out by Hall and Vinen at high temperatures, they are already too low at 1.3° . This tends to confirm the complicated character of the interaction of the rotons with the vortices, which is not reduced to simple scattering. When the rotons fly past at a small distance from the axis of the vortex processes apparently take place which have the character of "strong interaction". They lead to the transmission of an impulse of the order of the total roton impulse. Also the "smearing out" of the vortex because of its eigenoscillations may play an essential part. The attempt can be made to describe this interaction in a phenomenological manner by introducing a temperature-independent cross section ("breadth")

CARD 2/3

On the Absorption of the Second Sound in Rotating Helium II. 56-2-35/47

of the vortex, which corresponds to the transmission of the total roton impulse to the vortex. The determination of further experimental data would be desirable.

There is 1 figure and 1 Slavic reference.

ASSOCIATION: Institute for Physical Problems AN USSR.
SUBMITTED: (Institut fizicheskikh problem Akademii nauk SSSR)
MAY 10, 1957.
AVAILABLE: Library of Congress.

CARD 3/3

GOR'KOV, L.P.; PITAYEVSKIY, L.P.

Concerning scattering of light in He^3 and He^4 mixtures. Zhur. eksp.
i teor. fiz. 33 no.3:634-636 S '57. (MLRA 10:11)

1. Institut fizicheskikh problem AN SSSR.
(Light--Scattering) (Helium--Isotopes)

PITAYEVSKIY, L. P. and LIFSHITZ, E. P.

"The Theory of Superfluidity,"

report submitted but not presented at the International Union Conference on Low Temperature Physics, Leiden, 1951.

PITAYEVSKIY, L. P., Cand Phys-Math Sci -- (diss, "Studies
on the ~~Theory~~ Hyper-Viscosity of Liquid Helium. ~~Theory~~." Mos, 1958.
10 pp, (Acad Sci USSR, Inst ~~for~~ Phys Problems). 185 copies.
Bibliogr~~v~~ at ~~the~~ end of the text (12 titles), (KL 40-37,112)

AUTHOR: ~~Pilayarov, L. P.~~ 9-14-1-71, 11

TITLE: On the problem of the Anomalous Skin Effect in the Infrared Range (K voprosu ob anomal'nom skin-effekte v infrakrasnoy oblasti)

PERIODICAL: Zhurnal eksperimental'noy i teoreticheskoy fiziki, 1971, Vol. 54, Nr 4, pp. 942 - 946 (USSR)

ABSTRACT: The aim of this work is the consideration of the influence of the collisions between the electrons upon the optical properties of the metals in the infrared spectrum range; it computes the impedance of the metals at low temperatures. The frequency of the incident light is assumed to satisfy the conditions $\omega l/v \gg 1$ and $\omega \delta/v \gg 1$. v - denoting the velocity of the electrons, l - the free path, and δ the depth of penetration of the field into the metal. Furthermore the metal is assumed to have a high negative dielectric constant. The absorption is small and is considered as disturbance. This condition limits the usability of the formulae towards the side of the high frequencies. The computations were made for the case of very low temperatures, where $kT \ll \hbar\omega$ is valid. The

Card 1/3

On the Problem of the Anomalous Skin Effect in the
Infrared Range

50-14-4-1, 2

with various assumptions as to the law of the dispersion
of electrons, the form of the scattering probabilities, and on
the intensity of the interaction between the electrons. Above
all there is no reason for regarding the law of the dispersion
of the electrons to be constant square and their interaction to
be weak. Particular attention is given to the current of an excited elec-
tron or - strictly speaking - of a Fermi excitation - to be com-
puted. The author starts with the computation of the dielectric
constant. The so-called "number of free electrons" N depends
essentially on the interaction of the electrons. The electrons
in the metal can absorb the light in the collisions with the
atoms of the lattice (or with the lattice defects), in the
collisions with one another or with the surface of the metal,
or also in the absorption or emission of a photon.
In the case investigated the total absorption, clearly, is equal
to the sum of the terms which correspond to these processes.
For the energy balance, in case of isotropic symmetry is adopted
in the final result formula is given. Subsequently, the author
computes the collisions with the admixtures and with the surface
as well as the interelectronic collisions. The author thanks

Card 2, 3

On the problem of the wave-like Stark Effect in the Infrared Range

L. P. Lukatskiy, Member, Academy of Sciences, USSR, for Physics and Mathematics. There are no references, 7 of which are Soviet

ASSOCIATION Institute for the Study of Problems in Physics, AS USSR (Institute for the Study of Problems in Physics, AS USSR)

CLASSIFIED November 1, 1952

L. Math. --Optical properties

Card 3, 3

AUTHORS:

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..... (1958)

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..... must be attributed to the of the

..... of the liquid. ... is different from zero in H_1 and equal to zero in H_2 . Taking into account the quantum factor of the phenomena in liquid helium it is natural to choose as such a parameter the wave function $\psi(x,y,z) = \eta e^{i\phi}$ which plays the role of an "order parameter" of the superfluid part of the liquid. This paper deals only with those states of the normal part of which is assumed to be at rest. The velocity v of the superfluid part is zero. In this case the thermodynamic potential Ω is given by ... and the thermodynamic potential is given as

444

the Theory of Helium II

SOV/42-11-1-1-1

$$\int -1 \cdot \text{div} \mathbf{v} - \hbar^2 \text{grad}^2 \Psi \cdot \text{grad} \Psi / 2m \cdot \text{grad} \Psi$$
 with the boundary conditions $\Psi^* = \Psi$ and $\Psi = 0$ on the boundary. The equation $\Delta \Psi = 0$ is an analogon of the equation used in the hydrodynamical theory of superconduction is obtained. From the equation $-\hbar^2 \text{grad}^2 \Psi = 2m \cdot \text{grad} \Psi \cdot \text{grad} \Psi$ is obtained. To this equation belongs also the boundary condition $\Psi = 0$ which is to be used also for the free surface of helium. The thermodynamic potential Ω is expanded as a function of the order parameter into powers of $|\Psi|^2$

$$|\Psi|^2 = \frac{1}{2} \epsilon |\Psi|^2 + \frac{1}{4} \gamma |\Psi|^4$$

The theory used in this paper can be used only in the immediate neighborhood of the transition. The second section of the paper deals with some special problems. First the transition temperature and the properties of helium near a solid wall. In this case an additional surface energy appears. Then a helium film, i.e. a helium layer with the density $\rho = \rho_0$ is investigated. The temperature of the transition for a helium film is investigated in great masses of helium. Finally the properties of helium are investigated. There are figures and references. The figures are omitted.

307,

(Institute of Physics, AS UBR)

1. Helium (Liquid)—Properties
2. Helium (Liquid)—Mathematical analysis
3. Low temperature research

24 (0), 10 (4)

AUTHOR Pitayevskiy, L. P.

SOV/56-35-114/60

TITLE Phenomenological Theory of Superfluidity Near the λ -Point
(Fenomenologicheskaya teoriya sverkhtekuchesti vblizi
 λ -tochki)

PERIODICAL Zhurnal eksperimental'noy i teoreticheskoy fiziki, 1956,
Vol 35, Nr 2, pp 408-415 (USSR)

ABSTRACT: The present paper aims at generalizing the equation by Ginzburg and Pitayevskiy (Ref 1) for the non-steady case of superfluid helium II in the immediate neighborhood of the point of λ -transition. Proceeding from the energy- and momentum conservation theorem, the author derives the complete system of phenomenological equations which describe the behavior of superfluid helium at $T - T \ll T_\lambda$. The normal and the superfluid part are dealt with separately. The normal part is dealt with by the usual hydrodynamic method, and the superfluid part with the aid of an "effective wave function" formulated by the author. In conclusion, the author thanks L. D. Landau and V. L. Ginzburg for their valuable help and advice. There are 3 references, 2 of which are Soviet.

Card 1/1

Phenomenological Theory of Superfluidity Near
the λ -point

SOV/56-38-14/1

ASSOCIATION Institut fizicheskikh problem Akademii nauk SSSR
(Institute for Physical Problems, AS USSR)

SUBMITTED March 11, 1958

Card 1/2

BBV, 50-19-11-10-11

34, 27
AUTHOR: Pitajevskiy, L. P.
TITLE: Calculation of the Phonon Part of Mutual Frictional Force in Superfluid Helium (Vychisleniye fononnoy chasti sily vzaimnogo treniya v sverkhtekuchem gelii)
PERIODICAL: Zhurnal eksperimental'noy i teoreticheskoy fiziki, 1958, Vol 35, Nr 5, pp 1271-1275 (USSR)
ABSTRACT: Hall and Vinen (Kholi, Vaynen) (Ref 1) showed that in rotating helium II mutual frictional forces are active between the normal and the superfluid part of the liquid. For the purpose of calculating these forces it is necessary to investigate the scattering of elementary excitations consisting of the normal component of the liquid, i.e. of phonons and rotons, on the vortex lines. At temperatures $\lesssim 0.6^{\circ}\text{K}$ the phonon component in η_n is considerably smaller than that of rotons. In consequence, every mutual friction is dependent on the rotons at such temperatures. Calculation of this friction component was carried out by Hall and Vinen (Ref 2) as well as by Lifshits and Pitajevskiy (Ref 3). However, at such low temperatures the number of rotons decreases

Card 1/2

SOV 96-44-9-31, 32

Calculation of the Phonon Part of Mutual Frictional Force in Superfluid Helium

considerably, and the share contributed by phonons in mutual friction may be considerable. This contribution made by the phonon component for temperatures $T \lesssim 0.5^\circ\text{K}$ towards mutual friction is investigated by the present paper, i.e. the scattering of sound on vortex lines is calculated on the basis of ordinary hydrodynamical equations. It was found that the phonon component of mutual friction becomes comparable with the roton component. For the ratio between roton and phonon coefficients B_r and B_{ph} the following is obtained: $(B_{ph} \zeta_{n, ph} + B_r \zeta_{n, r}) / (\zeta_{n, ph} + \zeta_{n, r}) = B$ (total coefficient) $\zeta_n \ll \zeta_r$. In conclusion, the author thanks Professor Ye. M. Lifshits and L. D. Landau, Academicians, for their help and their discussions. There are 4 references, 2 of which are Soviet.

ASSOCIATION: Institut fizicheskikh problem Akademii nauk SSSR (Institute for Physical Problems of the Academy of Sciences, USSR)

SUBMITTED: June 13, 1958
Card 2/2

24(5)
AUTHOR:

U. M. KHARIN, L. I.

SOV. PHYS. JOURNAL
Spectrum

TITLE:

On the Properties of Elementary Excitation / Near the Disintegration
Threshold of the Excitations (O svoystvakh spetsialnykh elementarnykh
vozbuzhdeniy vbliz poroga raspada vozbuzhdeniy)

PERIODICAL:

Zhurnal eksperimental'noy i teoreticheskoy fiziki, 1976,
Vol 36, Nr 4, pp 1166-1176 (USSR)

ABSTRACT:

In the present paper the author investigates the properties of the excitation spectrum at its end for a Bose fluid. In the case of small momenta, excitation in a Bose liquid occurs in form of phonons, i.e. the excitation energy depends linearly on the momentum. With growing momentum, the spectrum more and more deviates from the linear spectrum, and finally its course depends on certain interaction properties of the particles of the fluid and can no longer be theoretically calculated in a general form. Whereas, e.g., the energy spectrum of liquid ⁴He at $p = 2 \cdot 10^{-19}$ g/cm/sec has a complex form, the excitation energy, with a further increase of momentum, attains a certain threshold value, above which excitation is unstable and may decompose into two or more excitations with lower energies.

Card 1/1

On the Properties of Elementary Excitation Spectrum SOV/58-104: 111
Near the Disintegration Threshold of the Excitations

This article on the continuous spectrum has interesting and original features which are investigated in the present paper. This investigation can be carried out in a general form by means of quantum field theoretical methods without it being necessary to assume weak interaction. First, the form of the spectrum near the threshold value is investigated. Three possible types of decay thresholds are found to exist. If the excitation energy is equal to the velocity of sound, the excitation leads to the production of phonons, whereas the results in excitation damping proportional to $(p - p_c)^{-1/2}$. At $p > p_c$ the excitation energy has a negative virtual contribution to $-av(\Delta p)^{1/2}$. An excitation in the center of mass is broken into two excitations with parallel momenta which are different from zero. The decomposition of excitation into two excitations with the momenta p which form an angle θ (cf graphs figure 1). In these two cases the spectral curves end at threshold point and the excitation velocity at this point is equal to that of one of the excitations produced in the decay. In the last part of the paper the author investigates the inelastic scattering

Card 2/3

On the Properties of Elementary Excitation Spectrum
Near the Disintegration Threshold of the Excitations

of neutrons. The case is investigated in which neutron energy losses are assumed to be $\xi \approx \xi_c$ and momentum losses $p \approx p_c$.

i.e. excitations with energies near decay threshold are investigated. All theoretical deliberations are based upon the laws of conservation. The author finally thanks Academician L. D. Landau for his advice, and V. M. Galitskiy, L. P. Gor'kov, and I. Ye. Dzyaloshinskiy for discussions. There are 3 figures and 5 Soviet references.

ASSOCIATION: Institut fizicheskikh problem Akademii nauk SSSR (Institute for Physical Problems of the Academy of Sciences, USSR)

SUBMITTED: October 6, 1958

Card 3/3

24(3)
AUTHORS: Dzyaloshinskiy, I. Ye., Pitayevskiy, L. P. SV, 56-76-t-25 10
TITLE: Van der Waals Forces in an Inhomogeneous Dielectric (Van-der-
Vaal'novy sily v neodnorodnom dielektrike)
PERIODICAL: Zhurnal eksperimental'noy i teoreticheskoy fiziki, 1959,
Vol 36, Nr 6, pp 1797-1805 (USSR)
ABSTRACT: In the present paper the forces are calculated which result
from the interaction of particles in an inhomogeneous
dielectric with the fluctuations of a long-wave electromagnetic
field (in which the dielectric is located). These forces may
be described as Van der Waal forces, because they are of a
similar nature. The contribution made by long-wave fluctuations
to the free energy is small compared to the total free energy
of the body, but it represents a new effect, namely that of
the non-additivity of the free energy of the body. The calcula-
tion of the (non-additive) correction to the free energy of
an inhomogeneous dielectric carried out by the authors cannot
be performed in the usual manner by determination of the field
energy in the medium because of the dissipation occurring in
the absorbing medium e.g. in the case of variable fields.
The authors used the diagram technique developed by

Card 1/2

Van der Waals Forces in an Inhomogeneous Dielectric

SOV/56-36-0-29/66

Matsubara (Ref 4), which was used also in the papers by Abrikosov, Gor'kov, Dzyaloshinskiy, and Fradkin (Ref 5). The authors first investigate the properties of the Green functions of an electromagnetic field in an inhomogeneous absorbing medium. In the second part of this paper the correction to the free energy of the system is derived by summation of the Matsubara graphs (Fig 2), and the corresponding part of the stress tensor (stress tensor of Van der Waal forces) is calculated. In an appendix the authors derive formulas for the Green functions in a homogeneous absorbing medium with complex dielectric constant. In conclusion, they thank Academician L. D. Landau and Ye. M. Lifshits for their interest and discussions. There are 2 figures and 8 references, 7 of which are Soviet.

ASSOCIATION: Institut fizicheskikh problem Akademii nauk SSSR (Institute for Physical Problems of the Academy of Sciences, USSR)

SUBMITTED: December 17, 1958

Card 2/2

10 (4)

AUTHORS: Dzyaloshinskiy, I. Ye., Lifshits, Ye. M., SOV/56-37-1-36/64
Pitayevskiy, L. P.

TITLE: Van der Waals' Forces in Liquid Films (Van-der-Vaal'sovy sily
v zhidkikh plenkakh)

PERIODICAL: Zhurnal eksperimental'noy i teoreticheskoy fiziki, 1959, Vol 37,
Nr 1(7), pp 227 - 241 (USSR)

ABSTRACT: The authors find general formulas for the determination of the thermodynamic quantities (chemical potential) of a liquid film, and they find the limiting laws for the dependence of the chemical potential on the thickness of the film. The difficulties in the generalization of the formulas derived for the vacuum in the case in which the interspace between the bodies is filled with any medium, are now eliminated because of the general formulas (Ref 2) already derived for that part of the thermodynamic quantities of any absorbing medium which is conditioned by the electromagnetic fluctuation field with the wave lengths $\lambda \gg a$ (a denoting the interatomic distances). This field corresponds to those forces which have the same nature as the van der Waals' forces between the single molecules at large distances. At first, the stress tensor in a stratified absorbent medium is

Card 1/3

Van der Waals' Forces in Liquid Films

SOV/56-37-1-36/64

calculated, and in the next part the forces of molecular interaction between solids are determined. In the case of a metallic "intermediate layer" between the bodies, the force of molecular attraction passes from the law l^{-3} at "small" distances to the law l^{-5} at "large" distances. The authors then investigate a liquid film on the surface of a solid body. This film is assumed to be applied to a wall vertically arranged in the field of gravity. $F(l) + \rho g z = \text{const}$ is the condition for the constancy of the chemical potential along the system, for $F(l)$ is the part of its chemical potential μ depending on the film thickness. Thus, $\mu = \mu_0 + F(l)$, μ_0 denoting the chemical potential of the "massive liquid". Further, $\mu(l) + \rho g z = 0$, the function $\mu(l)$ determining all thermodynamic properties of the film. The authors then investigate some typical cases which may be present according to the character of the function $\mu(l)$: (a) If $\mu(l)$ is a monotonely falling, everywhere positive function, the liquid does not moisten the solid surface, and no field is formed. (b) If $\mu(l)$ is a monotonely increasing, everywhere negative function, this usually corresponds to a liquid

Card 2/3

Van der Waals' Forces in Liquid Films

SOV/56-37-1-36/64

which completely moistens a solid surface. On a vertical wall, a film with a thickness tending to zero at $z \rightarrow \infty$ is particularly formed. This decrease in thickness takes place at first according to the law $l \sim z^{-1/4}$, then according to $z^{-1/3}$. Subsequently, the contribution to the chemical potential caused by forces of nonelectromagnetic origin is estimated. Finally, some films of liquid helium are specially investigated. The authors thank the Academician L. D. Landau for the discussion of the problems investigated here, and Professor B. V. Deryagin for the supply of his papers. There are 3 figures and 21 references, 10 of which are Soviet.

ASSOCIATION: Institut fizicheskikh problem Akademii nauk SSSR (Institute of Physical Problems of the Academy of Sciences, USSR)

SUBMITTED: February 12, 1959 (initially), and March 27, 1959 (after revision)

Card 3/3

SOV/56-37-2-48/56

4(6)
AUTHOR:

Fitajevskiy, L. P.

TITLE:

The Attraction of Small Particles Suspended in a Liquid Over Wide Distances

PERIODICAL:

Zhurnal eksperimental'noy i teoreticheskoy fiziki, 1959, Vol 37, Nr 2(4), pp 577-578 (USSR)

ABSTRACT:

The present "Letter to the Editor" is intended to derive a formula describing the interaction energy of particles suspended in a liquid, which is in interrelation with the Van der Waals forces. The distance separating the particles from one another should be large compared to their dimensions. In principle, this task may be solved on the basis of the general theory of Van der Waals forces in dielectrics. The expression for the interaction forces of an arbitrary body in a medium may be obtained by a simple transformation from the corresponding expression for the interaction forces in the vacuum. The expression for the additional pressure in a medium with the dielectric constant ϵ may be obtained from the expression for the pressure in the vacuum. By using the formulas by Landau and Lifshits for the interaction energy in the vacuum, the following formulas are

Card 1, 2

The Attraction of Small Particles Suspended in a Liquid Over Wide Distances

SOV, 56-37-2-48; 56

deduced for a medium containing N-particles per unit of volume:

$$U = \frac{3k}{16\pi^3 R^6} \int_0^\infty \left(\frac{\partial \epsilon(1f)}{\partial N} \right)^2 \frac{\partial f}{\epsilon^2(1f)} \quad \text{at } R \ll \lambda_0$$

$$U = \frac{23kc}{64\pi^3 R^7 \epsilon^{5/2}(0)} \left(\frac{\partial \epsilon(0)}{\partial N} \right)^2 \quad \text{at } R \gg \lambda_0$$

These formulas describe not only the interactions between macroscopic particles in suspension, but also e.g. of molecules in solution. There are 3 Soviet references.

ASSOCIATION: Institut zemnogo magnetizma, ionosfery i rasprostraneniya radiovoln Akademii nauk SSSR (Institute for Terrestrial Magnetism, Ionosphere, and the Propagation of Radiowaves of the Academy of Sciences, USSR)

SUBMITTED: May 15, 1959

Card 2/2

PITAYEVSKIY, L.P.

Superfluidity of liquid He³. Zhur.eksp.i teor.fiz. 37 no.6:1794-
1807 D '59. (MIRA 14:10)

1. Institut zemnogo magnetizma, ionosfery i rasporstraneniya
radiovoln AN SSSR.
(Superfluidity) (Liquid helium)

9.9845 (1538)

3037
S/570/60/000/017/007/012
E032/E114

AUTHORS: Gor'kov, L.P., Dzyaloshinskiy, I.Ye., and Pitayevskiy, L.P.

TITLE: Calculations of fluctuations in quantities described by transport equations

SOURCE: Akademiya nauk SSSR. Institut zemnogo magnetizma, ionosfery i rasprostraneniya radiovoln. Trudy, no.17(27). Moscow, 1960. Rasprostraneniye radiovoln i ionosfera. 203-207

TEXT: The authors discuss fluctuations in quantities which can be described by transport equations, e.g. the equations of Boltzmann, Fokker-Planck and Landau in the case of a Coulomb interaction between the particles. The knowledge of these fluctuations is essential to the theory of scattering of electromagnetic waves in rarefied gases and electron plasma. The method employed is analogous to that used by L.D. Landau and Ye.M. Lifshits (Ref.2: Electrodynamics of uniform media M., Gostekhizdat 1957, Ref.3: ZhETF v.32 618, 1957). It consists in the introduction into the transport equation of additional random

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terms whose correlations are then determined on the basis of the general theory of fluctuations. For example, the Boltzmann equation is modified to read

$$\frac{\partial \psi}{\partial t} + (\underline{v} \cdot \nabla) \psi = J + \gamma \quad (1)$$

where the collision integral J is given by

$$J = \iint w(p_1, p'_1, p, p') \{ n_0(p_1) \psi(p'_1) + n_0(p'_1) \psi(p_1) - n_0(p') \psi(p) - n_0(p) \psi(p') \} d^3 p_1 d^3 p'_1 d^3 p' \quad (2)$$

and γ is the "random" collision integral. The problem consists in the evaluation of the average of $\psi(p, r, t) \psi(p', r', t')$. It is shown that this average is in fact given by:

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$$\begin{aligned}
& \overline{y(p, r, t) y(p', r', t')} = 2b(r - r')b(t - t') \times \\
& \times \left\{ -n_0(p') \iint w(p', p_1, p_1, p_1') n_0(p_1) d^3 p_1 d^3 p_1' - \right. \\
& - n_0(p) \iint w(p', p_1, p_1, p_1') n_0(p_1') d^3 p_1 d^3 p_1' - \\
& + \delta(p - p') n_0(p) \iiint w(p_1, p_1', p, p_1) n_0(p_1) d^3 p_1 d^3 p_1' d^3 p_1 - \\
& \left. + n_0(p) n_0(p') \iint w(p_1, p_1', p, p') d^3 p_1 d^3 p_1' \right\} \quad (9)
\end{aligned}$$

which is equivalent to the results obtained by B.B. Kadomtsev (Ref. 5: ZhETF, v. 32, 943, 1957). It can be shown that the introduction of the "random" collision integral into Eq. (1) does not upset the conservation of the number of particles, energy and momentum. Another transport equation considered is the following

$$\frac{\partial v}{\partial t} + (\underline{v} \nabla) v = - \text{div } j \quad (10)$$

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where \mathbf{j} is the current density in the momentum space. Here it is convenient to introduce a "random" current \mathbf{v} so that

$$\frac{\partial \mathbf{j}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{j} = - \text{div} (\mathbf{j} \otimes \mathbf{y})$$

Expressions analogous to Eq (9) are then derived. An account of the general theory of fluctuations on which these calculations are based is given in "Statistical Physics" by L.D. Landau and Ye.M. Lifshits (Ref.4 1zd. 3 M., Gostekhizdat 1951). The method can be used for fluctuations in the equations for fermi and bose gases. A.A. Abrikosov and I.M. Khalatnikov (Ref.7. ZhETF v.34, 198-1958) have used it to study light scattering in liquid He³. Acknowledgments are expressed to L.D. Landau and Ye.M. Lifshits for discussions. S.M. Rytov and B.B. Kadomtsev are mentioned in connection with their contributions to the theory of fluctuations. There are 7 Soviet bibl. references.

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FITZGERALD

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AUTHOR

Matayevskiy, L. I.

TITLE

The Problem of the Shape of the Spectrum of Elementary Excitations of Liquid Helium II

PERIODICAL

Zhurnal eksperimental'noy i teoreticheskoy fiziki
Vol. 39, No. 1, pp. 17-17

TEXT Data on the shape of the energy spectrum of liquid helium II in the range of energy $\epsilon \approx \Delta = 17.3^\circ\text{K}$ were published in the paper of Ref. 1. These data had been obtained by a study of inelastic neutron scattering. It was found that the function $\epsilon(p)$ has a negative second derivative in this range. The authors of the above-mentioned publication believe that this is indicative of the existence of a second maximum of $\epsilon(p)$. The present "Letter to the Editor" however shows that this fact cannot be explained in this way. In the paper of Ref. 2 the writer has shown that the curve of the energy spectrum of liquid helium II usually does not exceed $\epsilon = \Delta$. This is due to the

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The problem of the shape of the spectrum of
Elementary Excitations of Liquid Helium II

By B. P. ...
R. D. B. ...

fact that an elementary excitation with $\epsilon \gg \Delta$ can split into two
excitations with $\epsilon = \Delta$ at an energy of $\epsilon = \Delta$
after which it is continuous since the spectrum ends at this point. Thus
the spectrum has the form $\epsilon(p) = \Delta \exp[-a/(p - p_0)]$, where p_0 is the
momentum for which $\epsilon = \Delta$, a , and Δ are constants. The amplitude $\epsilon(p)$
Fig. 1 shows the entire course of $\epsilon(p)$. The probability that a neutron
will produce an excitation with an energy ϵ tends to zero with

$\epsilon \rightarrow \Delta$ according to the law $w = A \frac{\Delta - \epsilon}{a} \ln \frac{a}{\Delta - \epsilon}$.

energy attained $\epsilon = 12.1$ K, w really amounts to 1 percent of the
probability of an excitation with $\epsilon = \Delta$. There are 1 figure and
references in Soviet and ...

ADDITIONAL: Institut zemnogo magnetizma, i radiofyzi-
rasprostraneniya radiovoln AN SSSR. Instytut fiziki
Terrestrial Magnetism, Ionosphere, and Propagation
Radio Waves of the USSR

Card 1

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The Problem of the Shape of the Spectrum of
Elementary Excitations of Liquid Helium II

S/O56/60, 039, 001
B006/B007

SUBMITTED: May 11, 1960

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86923

24.2120

S/056/60/011/011/011/011
B006/B077

AUTHOR: Pitayevskiy, L. P.

TITLE: Electric Forces in a Transparent Dispersive Medium

PERIODICAL: Zhurnal eksperimental'noy i teoreticheskoy fiziki, 1960,
Vol. 39, No. 5(11), pp. 1450 - 1458

TEXT: The present work deals with an examination of the form of the stress tensor of a variable electric field in a fluid transparent dispersive medium; the tensor is averaged over the oscillation period. In a dielectric of an isotropic fluid which is in a constant electric field the tensor assumes the form:

X

$$\sigma_{ik} = -P_0 \delta_{ik} - \frac{E^2}{8\pi} \left[\epsilon - \chi \left(\frac{\partial \epsilon}{\partial \nu} \right)_T \right] \delta_{ik} + \epsilon \frac{E_i E_k}{4\pi}$$

$P_0 = P_0(\nu, T)$ denotes the pressure in the fluid without an electric field, ν density, and ϵ the dielectric constant of the fluid. Now the author examines how this relation changes if a dispersion is assumed, that

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Electric Forces in a Transparent Dispersive Medium

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B006/BC77

is ϵ is considered a function of the field frequency. $\epsilon(\omega)$ is a real function (a typical example of such a medium is e.g. an electron plasma where ω is large as compared to the effective collision number). It is found that for σ_{ik} the same expression is obtained, but averaged, that is E^2 is substituted by E^2 and $E_i E_k$ by $E_i E_k$. The following expression is obtained for the energy flux: $\vec{S} = \frac{c}{4\pi} [\vec{E} \times \vec{H}] + \frac{E^2}{8\pi} (\omega \frac{\partial \epsilon}{\partial \omega} - c \frac{\partial \epsilon}{\partial c}) \vec{v}$, where \vec{v} is the flow rate of the medium. This expression is exact except for terms $\sim (v/c)^2$. The behavior of the dielectric constant of a medium with time dependent parameters is investigated: $\epsilon = \epsilon(\omega, \lambda)$, $\lambda = \lambda(t)$; it is found that the expression for ϵ contains an imaginary term (which is real in the case of time-independent parameters):

$$\epsilon(\omega, t) = \epsilon_0(\omega, t) + \frac{1}{2} \frac{\partial^2 \epsilon_0(\omega, t)}{\partial \omega \partial t} .$$

Finally, the case is examined where

the fluid is in a variable electric and in a strong constant magnetic field. The following expression is obtained:

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Electric Forces in a Transparent Dispersive Medium

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$$\begin{aligned} \sigma_{ik} = & -\left(\rho_{00} + \frac{H^2}{8\pi}\right) \delta_{ik} + \frac{1}{4\pi} H_i H_k + \frac{1}{16\pi} \left\{ E_{0i}^* E_{0m} \left(\rho \frac{\partial \epsilon_{im}}{\partial \rho} - \frac{\epsilon_{im} + \epsilon_{im}^*}{2} \right) \delta_{ik} \dots \right. \\ & + \frac{1}{16\pi} (\epsilon_1 E_{0i}^* E_{0k} + \epsilon_2 (E_0 H) (E_{0i}^* H_k + H_i E_{0k}^*) + \text{K. C.}) \dots \\ & \left. + \frac{H_i H_k}{8\pi} \left\{ \frac{\partial \epsilon_1}{\partial H^2} E_0 E_0^* + \frac{\partial \epsilon_2}{\partial H^2} |E_0 H|^2 + i \frac{\partial \epsilon_3}{\partial H^2} [E_0^* E_0] H \right\} \right. \end{aligned} \quad (46)$$

The author thanks L. D. Landau for his advice and discussion, V. L. Ginzburg and I. Ye. Dzyaloshinskiy for their discussions, and M. A. Leontovich for his suggestions concerning the last part of this work. There is 1 Soviet reference.

ASSOCIATION: Institut zemnogo magnetizma, ionosfery i rasprostraneniya radiovoln Akademii nauk SSSR (Institute of Terrestrial Magnetism, Ionosphere and Propagation of Radio Waves of the Academy of Sciences USSR)

SUBMITTED: July 27, 1960

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D22H/D304

AUTHOR:

Pitayevskiy, L. P.

TITLE:

The question of the disturbance induced in the plasma by a uniformly moving body

PERIODICAL:

Relativnyy zhurnal. Geofizika, no. 10, 1961, 14
abstract 1068 (Geomagnetizm i aeronomiya, 1, n. 10, 1961, 14-20)

TEXT: Expressions were found by means of the kinetic equation for the Fourier component of the electron density disturbance arising during the movement of a body with a velocity much greater than the thermal rate of ions in the plasma. It is suggested that the plasma is situated in a constant magnetic field in which the Larmor radius of ions is much greater than the body's dimensions. The effective area of dispersion of electromagnetic waves with a wavelength much greater than the body's size was found on the basis of these formulas. [Abstracter's note: Complete translation.]
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1106
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A006/A101

AUTHORS: Al'pert, Ya. L., Pitayevskiy, L. P.

TITLE: On the scattering of electromagnetic waves by perturbations caused by a rapidly moving body in a plasma

PERIODICAL: Geomagnetizm i aeronomiya, v. 1, no. 5, 1961, pp. 1-4

TEXT: Information is given on results of a theoretical analysis made of the effective scattering cross section of electromagnetic waves from the trail of a body moving rapidly through a magnetized plasma. The formulae, developed previously by Pitayevskiy, take into account the effect of the magnetic field, the number of ion collisions, the size of the body (especially those smaller than the wavelength) and to some extent the electrical field in the vicinity of the body. Due to the complexity of the formulae used, calculations were made with electronic computers for a number of concrete cases. The results obtained make it possible to obtain quantitative and qualitative concepts of the nature of scattering effect in a rarefied plasma (particularly in the ionosphere) for the case when the length of the particle path considerably exceeds the dimensions of a body moving at a velocity greater than the thermal velocity of ions $\sqrt{8kT/m}$ and considerably lower.

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On the scattering of...

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than the thermal velocity of electrons $\sqrt{2kT/m}$. It is shown that the cross section of scattering is a sharply directed multilobe function; its main maximum lies in the direction of the image "reflection" of the wave from the vector of the permanent magnetic field. The dependence of the cross section of scattering on the wavelength at different altitudes of the ionosphere is analyzed. In a number of cases the differential effective cross section attains its maximum value when the body moves in a direction forming a small angle with the parallel of latitude. The nature of the scattering is found to be a function of the geometry and physical dimensions of the body, and depends only slightly on its properties, shape and surface. There are 15 figures, 2 tables and 4 literature references.

ASSOCIATIONS: Institut Terrestrial Magnetism, Ionosphere and High-Frequency AN SSSR (Institute of Terrestrial Magnetism, Ionosphere and High-Frequency Radiation of Radiowaves, AN USSR); Institute of Geodesy and Aerial Photography I. Vavilova AN SSSR (Institute of Physics of the Academy of Sciences of the USSR, AN USSR)

SUBMITTED: August 4, 1966

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AUTHORS Prityayevskiy, L. I., Krestin, V. I.

TITLE Disturbances which occur when bodies are moving in a plasma

PERIODICAL Zhurnal teoreticheskoy fiziki, v. 40, no. 1, 1964, pp. 1-11

TEXT. A problem which lately has become typical is the scattering of electromagnetic waves by the track of a body moving in an isotropic electron plasma. The present paper deals with the theoretical study of this problem. The plasma is assumed to be infinite to such an extent that the mean free path of ions is large compared to both the length of the scattered electromagnetic wave and the dimensions of the body ($l \gg \lambda, R$). The scattering problem in question can be treated in two parts. Scattering by the body itself (e.g., a metal sphere) and scattering on a track formed by the sphere in the plasma; the latter is a series of disturbed electron concentrations. The scattering by the body itself can be described by conventional formulas of the diffraction theory and is not investigated here any further; however, it is shown that the part of a track of the body...

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same length. Scattering ... total scattering effect ... the perturbation ... the change of the dielectric constant of the plasma with a disturbed density ... is described by the relation

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$$\Delta \epsilon(\vec{k}) = \frac{4\pi e^2}{\omega^2} \int d^3r \rho(\vec{r}, t) e^{-i\vec{k}\cdot\vec{r}} \quad \text{and} \quad \rho(\vec{r}, t) = \frac{en_0}{4\pi R} \int d\Omega \int d\Omega' \vec{k} \cdot \vec{k}' E_{\vec{k}'} e^{i(\vec{k}-\vec{k}')\cdot\vec{r}} \quad \text{with } \vec{k}'$$

denoting the amplitude of the scattered wave \vec{k}' , the wave vector of the scattered wave $\vec{k}' = k' \hat{n}'$, the Fourier component of the disturbance of the electron density $\rho(\vec{r}, t)$ by

$$n_q = \int d\Omega \int d\Omega' \vec{k} \cdot \vec{k}' E_{\vec{k}'} e^{i(\vec{k}-\vec{k}')\cdot\vec{r}} \quad \text{where } \vec{k} \text{ is the wave vector of the incident wave, } \Psi \text{ the scattering angle between } \vec{k} \text{ and } \vec{k}'. \text{ The cross section in a solid angle } d\Omega \text{ is}$$

$$d\sigma = \frac{1}{k} \frac{d^3k'}{d\Omega'} \frac{1}{k'} \sin^2 \Psi \quad \text{It is not necessary to determine } d \text{ if it is necessary}$$

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to know $n_{\vec{q}}$. $n_{\vec{q}}$ may be determined by formulas of A. V. Gurevich, but here it is shown that it is easier to determine it directly from the equation of motion. The method brought here to determine $n_{\vec{q}}$ is also more exact, and it is possible to take into account effects occurring at small \vec{q} , which is not possible with the Gurevich method. General formulas are first derived for the case where the body is moving much slower than the thermal electrons ($v_0 \ll \sqrt{kT/m}$). In this case the electron density is a function of the potential according to Boltzmann: $n = n_0 \exp(e\psi/kT)$. After extensive calculations the following expressions are obtained:

$$n_{\vec{q}} = \frac{1}{iq} \frac{I(\vec{u})}{n(\vec{u}-\vec{v}_0) - 1\delta} d^3u / [2 - 2a \int_0^a e^{x^2} dx - i\sqrt{\pi/2} e^{-a^2}] ; a = n\vec{v}_0 \sqrt{M/2kT};$$

$\vec{u} = \vec{v} + \vec{v}_0$, \vec{v} is the ion velocity in a coordinate system moving with the body, M is the ion mass. The electron density decreases in proportion to $1/x^2$; this agrees with Gurevich. Furthermore, a formula is derived for $I(\vec{u})$; its calculation requires the knowledge of the law of ion scattering on a body with the electric field being taken into account. Though this

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formula does not permit the computation of $I(\vec{u})$ in the general case, it is possible to determine $I(\vec{u})$ for some special cases, as shown next. The calculation of the "collision integral" $I(\vec{u})$ is done for a) a slowly moving body ($V_0 \ll \sqrt{kT/M}$); b) a fast moving body ($V_0 \gg \sqrt{kT/M}$) with dimensions that are not small compared to the Debye radius; c) a small charged body ($eQ \ll R_D M v_1^2$). After this, $d\sigma$ and $n_{\vec{q}}$, respectively, is calculated for a slowly moving body, a fast moving large body

$$n_{\vec{q}} = \frac{n_0 q}{q} \left\{ \left[\frac{\pi}{2} - \sqrt{\pi} \left(\frac{M V_0^2}{2kT} \right)^{1/2} e^{-a^2} \right] + i 2 \left(\frac{M V_0^2}{2kT} \right)^{1/2} e^{-a^2} \int_0^a e^{x^2} dx \right\} \times \\ \times \left[2 \left(1 - a e^{-a^2} \int_0^a e^{x^2} dx \right) - i a \sqrt{\pi} e^{-a^2} \right]^{-1} \quad (41)$$

$$d\sigma = \frac{4\pi n_0^2 q^2}{16\pi^3} \left(\frac{a_0}{c} \right)^2 \frac{a_0^2}{q^2} \left\{ \left[\frac{\pi}{2} - \sqrt{\pi} \left(\frac{M V_0^2}{2kT} \right)^{1/2} e^{-a^2} \right]^2 + 4 \left(\frac{M V_0^2}{2kT} \right) \left(e^{-a^2} \int_0^a e^{x^2} dx \right)^2 \right\} \times \\ \times \left[4 \left(1 - a e^{-a^2} \int_0^a e^{x^2} dx \right)^2 + a^2 \pi e^{-2a^2} \right]^{-1} \quad (42)$$

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and a fast moving small body

$$n_q = \frac{2\pi Q^2 r^2 n_0}{q (2k^2 T^2 M V_0^2)^{1/2}} \ln \frac{R_D}{r_0} [V_0^2 - (V_0 n)^2] \left[(1 - 2a^2) \left(\sqrt{\pi} + 2i \int_0^a e^{x^2} dx \right) e^{-a^2} + 2ia \right] \times \int \cdot$$

$$\times \left[2 - 2a \left(\int_0^a e^{x^2} dx - i \frac{\sqrt{\pi}}{2} \right) e^{-a^2} \right]^{-1}. \quad (45)$$

The authors thank Ya. L. Al'pert and A. V. Gurevich for discussions.
A. G. Sitenko and S. N. Stepanov are mentioned. There are 8 references:
6 Soviet-bloc and 1 non-Soviet-bloc.

ASSOCIATION: Institut zemnogo magnetizma, ionosfery i rasprostraneniya
radiovoln Akademii nauk SSSR (Institute of Terrestrial
Magnetism, Ionosphere, and Propagation of Radiowaves, Academy
of Sciences USSR)

SUBMITTED: July 27, 1960

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16 7600

AUTHOR: Pitayevskiy, L. P.

TITLE: Vortex filaments in a rotating Bose gas

ABSTRACT: Zhurnal eksperimental'noy i teoreticheskoy fiziki, v. 4, no. 1, 1961, pp. 1-7

TEXT: According to Onsager (Navya 1949, 2, 449, 1949) and Feynman a superfluid liquid may contain vortex filaments, i.e., singular lines around which the superfluid part of the liquid rotates with the velocity $v = s\hbar/mr$ (m - mass of an atom of the liquid, r - radial distance from the vortex axis, s - an integer). The existence of these vortex filaments in rotating He II which has been proved experimentally by Hall and Vinen, is undoubted. The theoretical prediction of these vortex filaments has, however, been made on the basis of semiquantitative considerations; the wave function applied could not be derived from general principles. An attempt is made here to find a mathematical description of the processes leading to the formation of vortex filaments by using a simple model. The model chosen is an imperfect Bose gas with

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Vortex filaments in a nonideal ...

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weak pairwise repulsion of the atoms. This has been accurately studied by N. N. Bogolyubov. Bogolyubov's method is generalized and used for spatially inhomogeneous states, so as to study the vortex filaments in the Bose gas. The start is made from the Hamiltonian of the system in second-quantization representation [1]. It is assumed that a short-range potential is concerned here and that the gas is correspondingly dilute, so that the effective radius of the potential is small compared with the particle distance. We then obtain from

$$\hat{H} = \int \left\{ -\frac{\hbar^2}{2m} \psi^* \Delta \psi + \frac{1}{2} \int U(r-r_1) \psi^* \psi_1^* \psi_1 \psi d\tau_1 \right\} d\tau \quad (4)$$

$$\hat{H} = \int \left\{ -\frac{\hbar^2}{2m} \psi^* \Delta \psi + \frac{1}{2} g \psi^* \psi^* \psi \psi \right\} d\tau. \quad (5)$$

$$g = \int U(r) d\tau. \quad (5);$$

this Hamiltonian corresponds to an equation of motion for the

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Heisenberg operator ψ : $i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \Delta \psi + g \psi^2 \dots$. This operator is decomposed into two parts: $\psi = a_0 + \dots$, $\hbar \ll \hbar_0$; a_0 is a number,

and a small additional operator. One therefore obtains $i\hbar \frac{\partial a_0}{\partial t} = -\frac{\hbar^2}{2m} \Delta a_0 + g a_0^2 a_0$. If $a_0 \sim e^{-iE_0 t/\hbar}$, one obtains

$-\frac{\hbar^2}{2m} \Delta a_0 - E_0 a_0 + g a_0^2 a_0 = 0$. An equation of this kind has already

been examined by V. L. Ginzburg and Litayevskiy (ZhETF, 34, 1959, 19, 6) in connection with the phenomenological theory of superfluidity near the λ point. The vortex filaments correspond to a solution of this equation, which is symmetrical with respect to a given axis. If one puts

$a_0 = e^{i\theta} F(r)$ and passes over to cylindrical coordinates, one obtains

$$-\frac{\hbar^2}{2m} \frac{1}{r} \frac{d}{dr} r \frac{dF}{dr} + \frac{\hbar^2}{2mr^2} F - E_0 F + gF^3 = 0. \quad (10)$$

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Vortex filaments in a nonideal ...

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(10) has a solution vanishing with $r \rightarrow 0$, and tending to a constant, n_0 , with $r \rightarrow \infty$; n_0 corresponds to the particle density at r_∞ and coincides with the total gas density except for the terms $\sim g^{3/2}$. $E_0 = gn_0$, and the vortex filament at $s = 0$ is given by

$a_0 = \bar{n}_0 \exp[i(\omega - gn_0 t/\lambda)] \psi_0(r/r_0)$, $r = \lambda / 2\pi gn_0$, $\psi_0(\xi)$ is a real function, for $\xi \rightarrow 0$ $\psi_0 \sim \xi^2$, and for $\xi \rightarrow \infty$ $\psi_0(\xi) \sim 1 - \xi^{-2}/2$. r_0 denotes the thickness of the vortex axis. The mean rotational velocity of the particles about the vortex filament is correctly found to be λ/mr . The approximation made with $v \ll a_0$ is justified if the scattering amplitude of the particles is small compared with their distance: $mg/\lambda^2 \ll n_0^{-1/3} \approx n^{-1/3}$, n - particle density; therefrom one has $r_0 \gg n^{-1/3}$, which constitutes a necessary condition. In what follows, the oscillation of a vortex line is examined and a branch with a dispersion

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0.056, 0.1, 0.15, 0.2, 0.25, 0.3, 0.35, 0.4, 0.45, 0.5, 0.55, 0.6, 0.65, 0.7, 0.75, 0.8, 0.85, 0.9, 0.95, 1.0

law is shown to appear in the energy spectrum of a Bose gas if there is a vortex filament. The energy spectrum of weakly excited states is determined from

$$i\hbar \frac{\partial \theta}{\partial t} = \frac{\hbar^2}{2m} \Delta \theta - 2g |a_0|^2 \theta - g_2 |\theta|^2 \theta, \quad (10)$$

with

$$\theta = e^{i(E_0 t - k r)} \sum_k b_k(k, r) e^{i(k_2 + l_2 \theta)} \quad (l = 0, \pm 1, \pm 2, \dots)$$

A system of linear equations

Vortex filaments in a nonideal ...

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$$i\hbar \frac{db_l(k)}{dt} = - \frac{\hbar^2}{2m} \left[\frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} - \frac{(l-1)^2}{r^2} - k^2 \right] b_l(k) - gn_0(2\psi_0^2 - 1)b_l(k) - gn_0\psi_0^2 b_{-l}^*(-k)$$

$$i\hbar \frac{db_{-l}^*(-k)}{dt} = - \frac{\hbar^2}{2m} \left[\frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} - \frac{(l-1)^2}{r^2} - k^2 \right] b_{-l}^*(-k) - gn_0(2\psi_0^2 - 1)b_{-l}^*(-k) - gn_0\psi_0^2 b_l(k)$$

is obtained for $b_l(k)$ and $b_{-l}^*(-k)$. It is shown that among the solutions of (20) there is actually one such that corresponds to longwave oscillations (corresponding to macroscopic vortices with thick vortex filaments in the diluted gas). The dispersion law:

$$\omega = \frac{\hbar k^2}{2m} \ln \frac{1}{kr_0} = \frac{\hbar k^2}{2m} \ln \frac{\sqrt{2gm n_0}}{k\lambda}$$

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Vortex filaments in a plasma.

1966, 61, 4, 1013-1017
Phys. Rev.

is obtained, which corresponds to that of an ideal fluid. Academician L. D. Landau is thanked for his discussions. There are 11 references, 4 Soviet-bloc and 7 non-Soviet-bloc.

ASSOCIATION: Institut zemnogo magnetizma, ionosfery i rasprostraneniya radiovoln Akademii nauk SSSR (Institute of Terrestrial Magnetism, Ionosphere, and Propagation of Radio Waves, Academy of Sciences, USSR)

SUBMITTED: September 14, 1966

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S/053/61/073/003/001/004

B125/B201

24.6100

AUTHORS: Dzyaloshinskiy, I. Ye., Lifshits, Ye. M., and
Pitayevskiy, L. P.

TITLE: General theory of the Van der Waals forces

PERIODICAL: Uspekhi fizicheskikh nauk, v. 73, no. 3, 1961, 381-422

TEXT: A brief report is first given of the methods of the quantum field theory, and the general theory of the Van der Waals forces is then explained on this basis. Such a theory has been developed for the first time by Ye. M. Lifshits. The application of the methods of the quantum field theory to the problems of statistical physics at finite temperature is based on a paper by Matsubara. According to it, the free energy can be calculated by the rules of Feynman's graph technique. Matsubara's technique can be appreciably improved by taking account of some general properties of the Green functions (A. A. Abrikosov, L. P. Gor'kov, I. Ye. Dzyaloshinskiy, ZhETF 33, 799 (1959), Ye. S. Fradkin, ZhETF 36, 1286 (1959). The following series presented schematically must be summed

X

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General theory of the...

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to calculate the total Green function of the photon:

$$D = \text{---} + \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} \text{---} + \dots$$

In the case of a spatially inhomogeneous system it has the form

$$D_{\alpha\beta}(r_1, r_2; \xi_n) = D_{\alpha\beta}^0(r_1, r_2; \xi_n) + \int D_{\alpha\gamma}(r_1, r_3; \xi_n) \Pi_{\gamma\delta}(r_3, r_4; \xi_n) \times \times D_{\delta\beta}^0(r_4, r_2; \xi_n) dr_3 dr_4 + \int D_{\alpha\gamma}(r_1, r_3; \xi_n) \Pi_{\gamma\delta}(r_3, r_4; \xi_n) D_{\delta\mu}^0(r_4, r_5; \xi_n) \times \times \Pi_{\mu\nu}(r_5, r_6; \xi_n) D_{\nu\beta}^0(r_6, r_2; \xi_n) dr_3 dr_4 dr_5 dr_6 + \dots \quad \text{Eq. 2.8} \quad (2.8)$$

$\Pi_{\alpha\beta}(r_1, r_2; \xi_n)$ signifies the so-called polarization operator of the system. (2.9) or, in another form,

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